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## A Hybrid Model for Multiscale Laser Plasma Simulations with Detailed Collisional Physics

AFOSR Plasma and Electroenergetics Review Meeting 29-30 November 2016



#### **Collaborators:**

- R. Caflisch (UCLA)
- A. Karagozian (UCLA)

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## **Outline**



- Goals
- Review of Past Work
- Argon Collisional-Radiative Complexity Reduction Validation
- Non-Maxwellian CR
- Phase-accurate Multiscale Particle Push



## Goals



**FRC** 

 Utilize hybridization techniques to produce accurate and efficient plasma simulations that spans many orders of magnitude in both space and time.

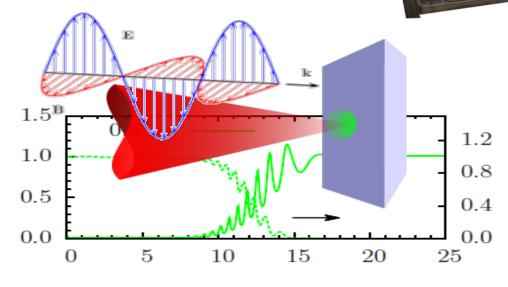
• Capture complex physics: excitation/ionization, transport, radiation, etc.

Consistent collision operator across different levels of fidelity.

## **Current Focus:**

- Generalization of collisionalradiative kinetics with level grouping
- General Hybridization techniques
- Focus on each solver before hybridization
- Special attention to low density low energy conditions







## **RQRS M&S Group**



#### Government

- Dr. David Bilyeu (LPI, in-space Chem)
- Dr. Justin Koo (Flight Support, Group Lead)
- Dr. Rob Martin (FRC)

## Onsite Contractors (ERC inc.)

- Richard Abrantes (Grad Student, LPI)
- Dr. Jun Araki (Flight Support, PIC)
- Dr. Carl Lederman (LPI)
- Dr. Michelle Scharfe (Flight Support)
- Dr. Eder Sousa (FRC)
- Jonathan Tran (Grad Student, Implicit PIC)

#### Summer Students

- Astrid Raisanen
  - PhD UofM; Vlasov HET
- Daniel Crews
  - M.S. Washington; Collisionless Shock for V&V
- Kari Kawashima
  - B.S. UCLA; SM/MURF beta tester

#### Previous

- Dr. Hai Le (LPI; now at Livermore)
- Dr. Artan Qerushi (FRC now at Lockheed)



## **Summary of Past Work**



#### **Maxwellian Inelastic Collisions**

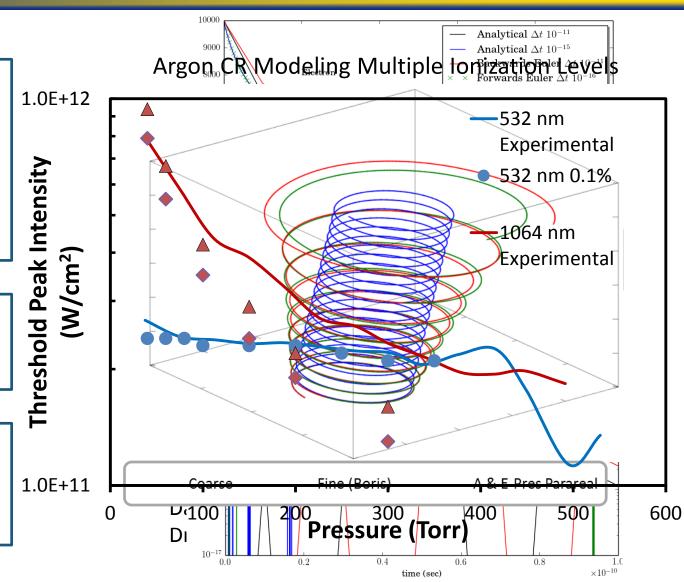
- Detailed CR model for multiple ionization stages
- Validation against experimental data
- Nonequilibrium radiation transport: coupling with a collisional-radiative model
- Inelastic collisions in a MF plasma: enhanced thermochemical kinetics.

#### Multiscale Hybridization

A time-parallel/multiscale method with energy preservation

#### Analytical BGK

- Conservative even for large mass ratios
- Conservation are independent of collisional frequencies





## **Collisional Radiative (CR) Overview**



#### Updates

- Expanded complexity reduction to include multiple ionization levels
- Adaptive integration technique -> fixes rate calculation for higher electron temperature
- Investigate grouping sensitivity
- Linked with LANL database for Argon cross sections and atomic level information
- Algorithms not hard coded for Argon.

#### **Levels of Complexity**

- Full rates (LANL)
- Cutoff nearly ionized levels
- Grouping Strategies
  - Uniform
  - Boltzmann
  - QSS (Boltzmann and Planck equilibrium)
- Group Selection
  - Electron configuration (no splitting information)
  - Highly excited states
  - Analysis of full run
  - Numerical optimization



## **CR Governing Equations**



$$\left(\frac{df_e(\varepsilon,t)}{dt}\right)_{\text{coll,ex/dex}} = -\sum_{m>n} N_n(f_e(\varepsilon_0)) \left(\sqrt{\frac{2\varepsilon_0}{m_e}}\right) \sigma_{(m|n)}^{e,\uparrow}(\varepsilon_0) 
+ \sum_{m>n} N_m(f_e(\varepsilon_1)) \left(\sqrt{\frac{2\varepsilon_1}{m_e}}\right) \sigma_{(n|m)}^{e,\downarrow}(\varepsilon_1) 
+ \sum_{mn} N_n(f_e(\varepsilon_0)) \left(\sqrt{\frac{2\varepsilon_0}{m_e}}\right) \sigma_{(m|n)}^{e,\downarrow}(\varepsilon_0)$$

$$\left(\frac{dN_{n}(t)}{dt}\right)_{\text{coll,ex/dex}} = -\sum_{m>n} N_{n} \int_{\varepsilon_{0}} (f_{e}(\varepsilon_{0})) \left(\sqrt{\frac{2\varepsilon_{0}}{m_{e}}}\right) \sigma_{(m|n)}^{e,\uparrow}(\varepsilon_{0}) d\varepsilon_{0} 
+ \sum_{m>n} N_{m} \int_{\varepsilon_{1}} (f_{e}(\varepsilon_{1})) \left(\sqrt{\frac{2\varepsilon_{1}}{m_{e}}}\right) \sigma_{(n|m)}^{e,\downarrow}(\varepsilon_{1}) d\varepsilon_{1} 
+ \sum_{m$$

$$\left(\frac{dN_n(t)}{dt}\right)_{\text{coll,i/r}} = -N_n \int_{\varepsilon_0} (f_e(\varepsilon_0)) \left(\sqrt{\frac{2\varepsilon_0}{m_e}}\right) \sigma_{(+|n)}^{e,\uparrow}(\varepsilon_0) d\varepsilon_0 + N_+ \int_{\varepsilon_0} \int_{\varepsilon_1} \int_{\varepsilon_2} (f_e(\varepsilon_1)) (f_e(\varepsilon_2)) \left(\sqrt{\frac{2\varepsilon_1}{m_e}}\right) \left(\sqrt{\frac{2\varepsilon_2}{m_e}}\right) \left(\frac{\partial^2}{\partial \varepsilon_1 \partial \varepsilon_2} \sigma_{(n|+)}^{e,\downarrow}(\varepsilon_1, \varepsilon_2; \varepsilon_0)\right) d\varepsilon_1 d\varepsilon_2$$



## **CR Governing Equations cont.**



#### **Current Assumptions**

- Electron dominated collisions
- $\delta$  function for ion distribution
- Maxwellian electrons

#### **Current Model Includes**

- Multiple ionization levels
- Excitation/de-excitation
- Ionization/recombination
- Multi-photon ionization and inverse Bermsstrahlung

Radiation losses via, Bound-**Bound and Bound-Free** 

$$\begin{split} \frac{dn_{n}^{+k}}{dt} &= \\ &- \sum_{m>n} \alpha_{(m|n)}^{+k,e} N_{e} N_{+k,n} + \sum_{m>n} \beta_{(n|m)}^{+k,e} N_{e} N_{+k,m} + \sum_{m>n} A_{(n|m)}^{+k} N_{+k,m} \\ &+ \sum_{m$$

\*previous included



## **CR Grouping Techniques**



## Uniform

- Solely conserves number density
- Weights each level according to level degeneracy within group

Conserved Variable:

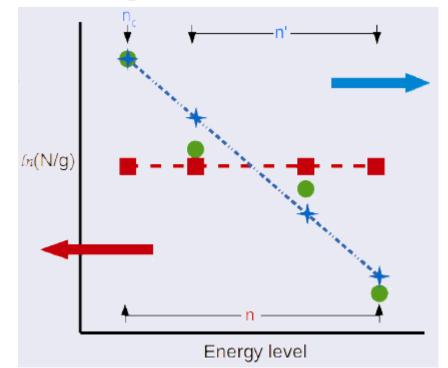
$$\widetilde{N}_n = \sum_{i \in n} N_i$$

Effective rate:

$$\tilde{\alpha}_{(m|n)} = \sum_{i \in n} \frac{g_i}{g_n} \sum_{j \in m} \alpha_{(j|i)}$$

$$\frac{dN_n}{dt} = N_e \left[ \sum_{m>n} \alpha_{(m|n)} N_n + \sum_{m< n} \beta_{(m|n)} N_n \right] + \cdots$$

$$\frac{d\widetilde{N}_n}{dt} = N_e \widetilde{N}_n \left[ \sum_{m > n} \widetilde{\alpha}_{(m|n)} + \sum_{m < n} \widetilde{\beta}_{(m|n)} \right] + \cdots$$



## **Boltzmann**

- Conserves number density
- Preserves energy in groups through group temperature description

Conserved Variable:

$$N_{n_0} \& N'_n = \frac{N_{n_0}}{g_{n_0}} \sum_{i \in n} g_i e_i^{-\Delta E_i/T_n}$$

Effective rate:

$$\tilde{\alpha}(m'|n') = \sum_{i \in \mathbf{n}} \frac{g_i e^{-\Delta E_i/T_n}}{Z'_n} \sum_{j \in \mathbf{m}'} \alpha_{(j|i)}$$



## **Complexity Reduction for Argon**



### **Full Lines (LANL)**

 Based on quantum calculations with corrections for low temperature

#### **NIST Cutoff**

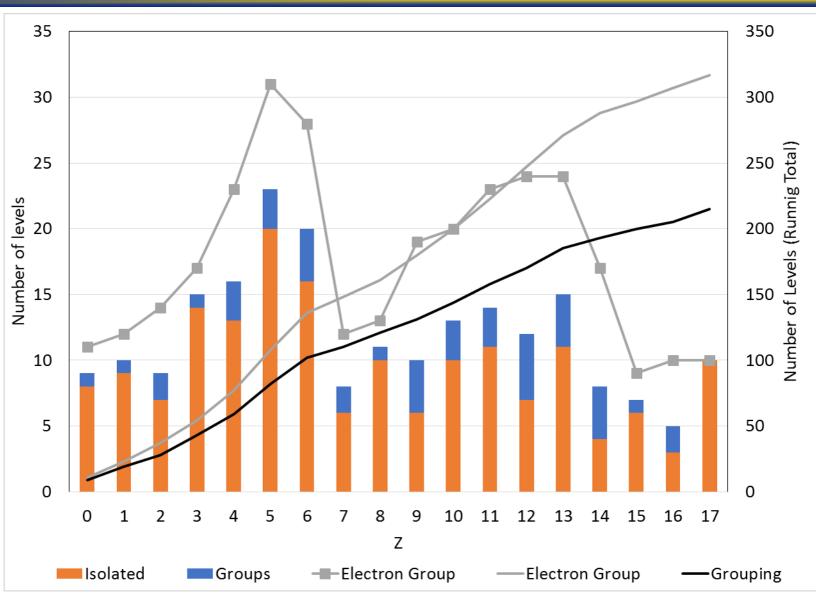
- Starts with LANL and assumes higher excited states are ionized
- Cutoff experimentally determined
- 2-3x reduction

### **Electron Configuration**

- Groups based on electron configuration
- Uses uniform grouping
- 10-15x reduction over NIST

### Grouping

- Boltzmann or Uniform grouping
- Saves 20-30% over Electron Splitting
- Case by case basis





# Argon Level Grouping Isothermal Test Setup



# Simulation Setup

- Pressure: 4.22 Torr (5.55x10<sup>-3</sup> atm)
- Ion Temperature: 0.035 eV
- Atomic Density: 10<sup>20</sup> 1/m<sup>3</sup>
- Ionization fraction: 10<sup>-13</sup>
- Electron Temperature: 10 & 100 eV
- $t = [0,10^6]$  seconds

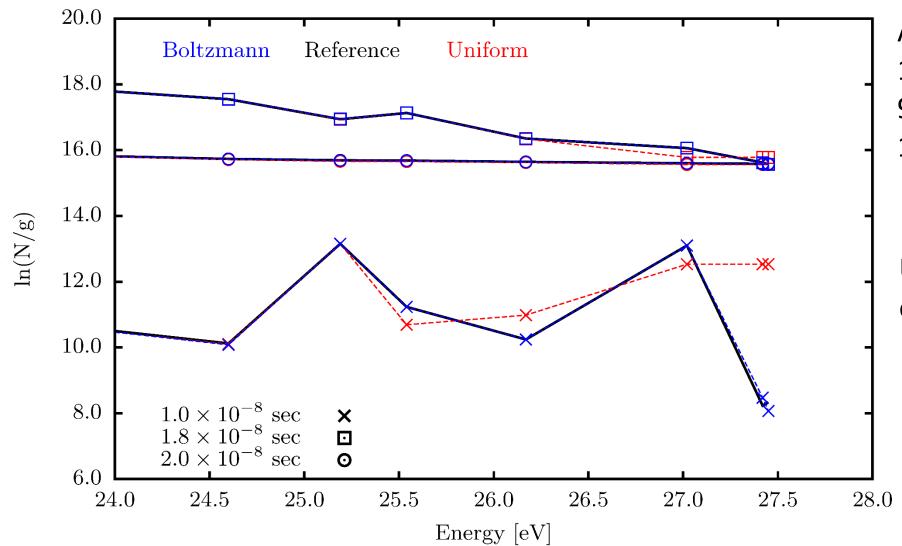
## Groupings

- NIST cutoff with electron grouping
- NIST cutoff with electron grouping and Boltzmann grouping
- NIST cutoff with electron grouping and Uniform grouping



# **Argon Level Grouping Isothermal Test Electron Temperature 10 eV**





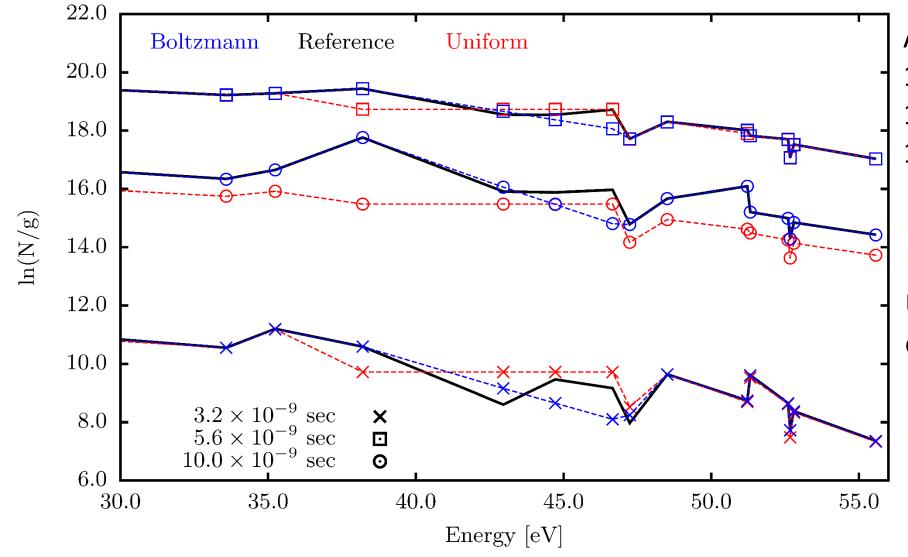
Argon +1
12 electron configurations
9 Isolates states
1 group with 3 states

Uniform group propagates errors to isolated states



# **Argon Level Grouping Isothermal Test Electron Temperature 100 eV**





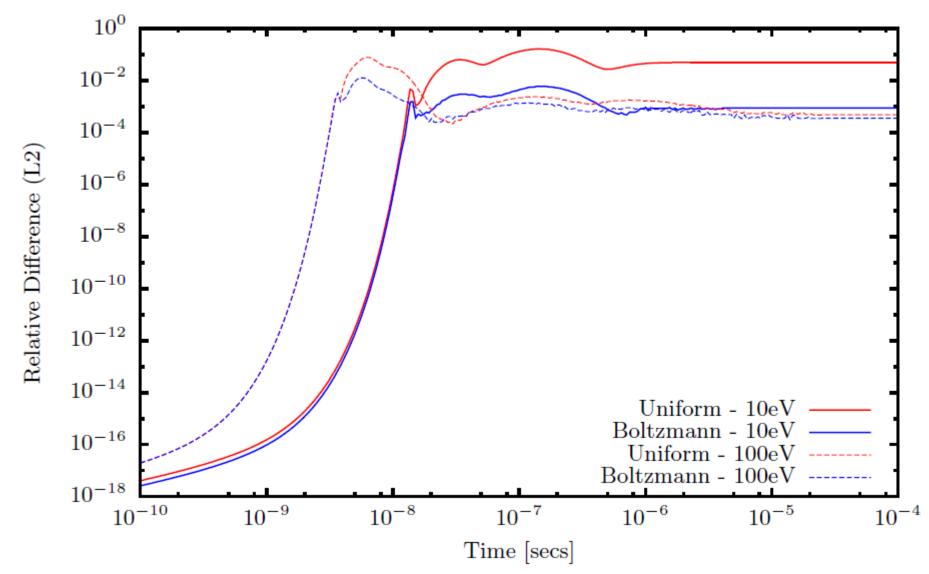
Argon +3
17 electron configurations
14 Isolates states
1 group with 3 states

Uniform groups propagates errors to isolated states



# Argon Level Grouping Isothermal Test Error





ODE Solver tolerance
Relative error 1<sup>-4</sup>
Absolute error 1
~1200 Radau5 time steps
dt increases exponentially

Solution improves if either tolerance is decreased but at the expense of computational time. E.g. relative error  $10^{-6}$  -> computational time triples



# Non-Maxwellian Electrons CR - Ions (Preliminary)



## Current Assumptions

- Electron dominated collisions
- Single ionization level
- Isotropic EEDF
- δ function for ion distribution

#### Current Model Includes

- Elastic electron collisions
- Excitation/de-excitation
- Ionization/recombination

$$\left(\frac{dN_{n}(t)}{dt}\right)_{\text{coll,ex/dex}} = -\sum_{m>n} N_{n} \int_{\varepsilon_{0}} (f_{e}(\varepsilon_{0})) \left(\sqrt{\frac{2\varepsilon_{0}}{m_{e}}}\right) \sigma_{(m|n)}^{e,\uparrow}(\varepsilon_{0}) d\varepsilon_{0} 
+ \sum_{m>n} N_{m} \int_{\varepsilon_{1}} (f_{e}(\varepsilon_{1})) \left(\sqrt{\frac{2\varepsilon_{1}}{m_{e}}}\right) \sigma_{(n|m)}^{e,\downarrow}(\varepsilon_{1}) d\varepsilon_{1} 
+ \sum_{m$$

$$\left(\frac{dN_n(t)}{dt}\right)_{\text{coll,i/r}} = -N_n \int_{\varepsilon_0} \left(f_e(\varepsilon_0)\right) \left(\sqrt{\frac{2\varepsilon_0}{m_e}}\right) \sigma_{(+|n)}^{e,\uparrow}(\varepsilon_0) d\varepsilon_0$$

$$-\int_{\varepsilon_0} \int_{\varepsilon_1} \int_{\varepsilon_2} \left(f_e(\varepsilon_1)\right) \left(f_e(\varepsilon_2)\right) \left(\sqrt{\frac{2\varepsilon_1}{m_e}}\right) \left(\sqrt{\frac{2\varepsilon_2}{m_e}}\right) \left(\frac{\partial^2}{\partial \varepsilon_1 \partial \varepsilon_2} \sigma_{(n|+)}^{e,\downarrow}(\varepsilon_1, \varepsilon_2; \varepsilon_0)\right) d\varepsilon_1 d\varepsilon_2$$



## **Non-Maxwellian Electrons CR - Electrons**



## (Preliminary)

## Current Assumptions

- Electron dominated collisions
- Single ionization level
- Isotropic EEDF
- δ function for ion distribution

### Current Model Includes

- Elastic electron collisions
- Excitation/de-excitation
- Ionization/recombination

$$\begin{pmatrix} \frac{df_{e}(\varepsilon,t)}{dt} \end{pmatrix}_{\text{coll,ex/dex}} = -\sum_{m>n} N_{n}(f_{e}(\varepsilon_{0})) \left(\sqrt{\frac{2\varepsilon_{0}}{m_{e}}}\right) \sigma_{(m|n)}^{e,\uparrow}(\varepsilon_{0}) \\ + \sum_{m>n} N_{m}(f_{e}(\varepsilon_{1})) \left(\sqrt{\frac{2\varepsilon_{1}}{m_{e}}}\right) \sigma_{(n|m)}^{e,\downarrow}(\varepsilon_{1}) \\ + \sum_{mn} N_{n}(f_{e}(\varepsilon_{0})) \left(\sqrt{\frac{2\varepsilon_{0}}{m_{e}}}\right) \sigma_{(m|n)}^{e,\downarrow}(\varepsilon_{0}) \\ \begin{pmatrix} \frac{df_{e}(\varepsilon,t)}{dt} \\ coll,i/r \end{pmatrix} = \\ - \sum_{n} N_{n}(f_{e}(\varepsilon)) \left(\sqrt{\frac{2\varepsilon_{0}}{m_{e}}}\right) \sigma_{(+|n)}^{e,\uparrow}(\varepsilon_{0}) \\ + N_{+} \int_{\varepsilon_{1}} \int_{\varepsilon_{2}} (f_{e}(\varepsilon_{1})) (f_{e}(\varepsilon_{2})) \left(\sqrt{\frac{2\varepsilon_{1}}{m_{e}}}\right) \left(\sqrt{\frac{2\varepsilon_{2}}{m_{e}}}\right) \left(\frac{\partial^{2}}{\partial \varepsilon_{1}\partial \varepsilon_{2}} \sigma_{(n|+)}^{e,\downarrow}(\varepsilon_{1},\varepsilon_{2};\varepsilon_{0})\right) d\varepsilon_{1} d\varepsilon_{2} \\ - N_{+} (f_{e}(\varepsilon_{2})) \left(\sqrt{\frac{2\varepsilon_{2}}{m_{e}}}\right) \int_{\varepsilon_{0}} \int_{\varepsilon_{1}} (f_{e}(\varepsilon_{1})) \left(\sqrt{\frac{2\varepsilon_{1}}{m_{e}}}\right) \left(\frac{\partial^{2}}{\partial \varepsilon_{1}\partial \varepsilon_{2}} \sigma_{(n|+)}^{e,\downarrow}(\varepsilon_{1},\varepsilon_{2};\varepsilon_{0})\right) d\varepsilon_{1} d\varepsilon_{0} \\ + \sum_{n} N_{n} \int_{\varepsilon_{1}} \int_{\varepsilon_{0}} \left(\sqrt{\frac{2\varepsilon_{0}}{m_{e}}}\right) (f_{e}(\varepsilon_{0})) \left(\frac{\partial^{2}}{\partial \varepsilon_{1}\partial \varepsilon_{2}} \sigma_{(+|n)}^{e,\downarrow}(\varepsilon_{2};\varepsilon_{0},\varepsilon_{1})\right) d\varepsilon_{0} d\varepsilon_{1} d\varepsilon_{0} \\ \end{pmatrix}$$

 $= \nu_{ee}(F_e - f_e) + \nu_{ei}(F_{ei} - f_e)$ 



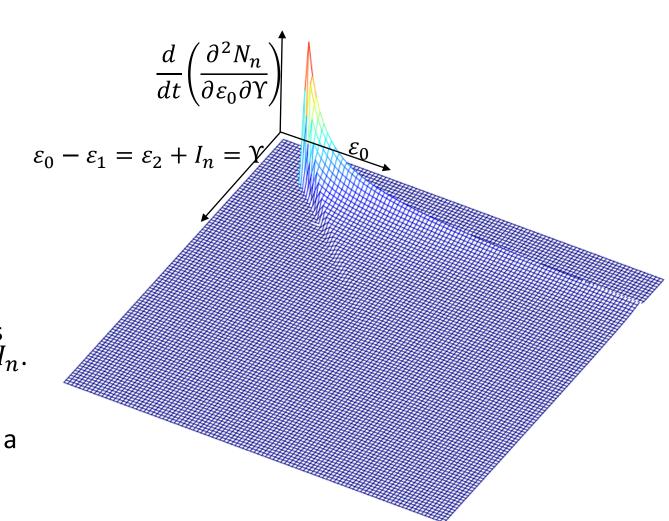
# Non-Maxwellian CR Numerical Challenges



- Long complex CR formulas
- Stiff equations
- Range of scales
- Boundary conditions
- Multi-dimensional integrations
- Hydrogen recombination example:

• 
$$e_0 + H_n \leftarrow e_1 + e_2 + H_+$$

- Energy equation in terms of the electron's kinetic energies,  $\varepsilon$ , and ionization energy  $I_n$ .
- $\varepsilon_0 = \varepsilon_1 + \varepsilon_2 + I_n$
- Evaluating the effect of recombination on a single species of hydrogen requires the evaluation of a 2D integral.





# **Hybridization Techniques**



- Solve ODE/PDE with a coarse "C" solver and fine "F" solver
- Coarse propagates less information and is more computationally efficient
- Fine propagates more information and is more accurate
- User defined coarse error function h(u),  $0 \le h \le 1$
- when h=0 coarse is accurate
- Q compression operator
- R reconstruction operator
- h, ε, Q, & R, are problem dependent

Simpler hybrid (H) method

$$H1_{[t+dt,t]} = \begin{cases} C_{[t+dt,t]} \big( Q \text{ (if needed)} \big) & \text{if } h \leq \epsilon \\ F_{[t+dt,t]} \big( R \text{ (if needed)} \big) & \text{if } h > \epsilon \end{cases}$$

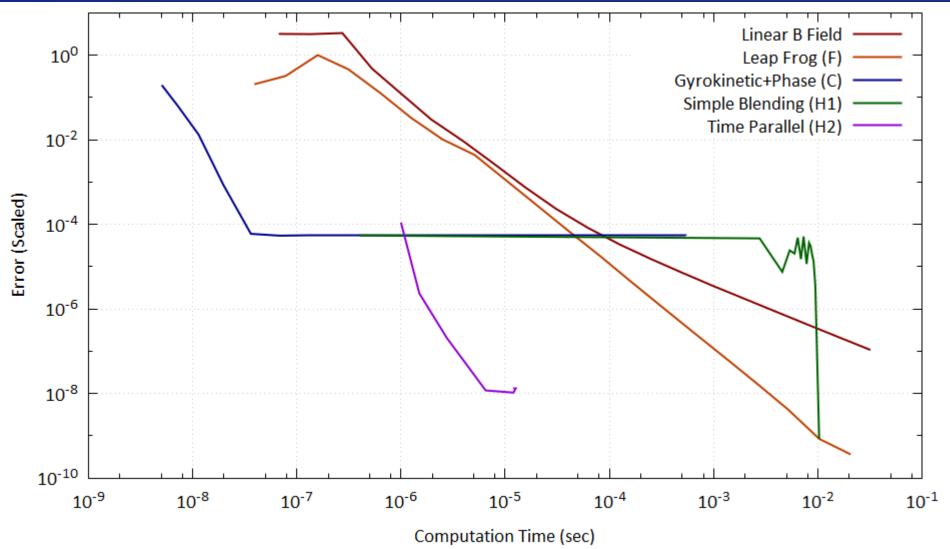
More complex time parallel (TP) method

$$H2_{[Jdt,0]}(u(0)) = \begin{cases} C_{[Jdt,0]}Q(u(0)) & \text{if } h \leq \epsilon \\ C_{[Jdt,0]}Q(u(0)) + \\ \sum_{j=0}^{j=J-1} RC_{[Jdt,(j+1)dt]}QF_{(j+1)dt,jdt]}RC_{[jdt,0]}Q(u(0)) \\ \sum_{j=0}^{j=J-1} RC_{[Jdt,(j+1)dt]}C_{(j+1)dt,jdt]}C_{[jdt,0]}Q(u(0)) \end{cases} \text{ if } \epsilon < h \leq \gamma$$
 and 
$$F_{[Jdt,0]}(u(0)) & \text{if } h > \gamma$$



## **Hybridization Charged Particle in Magnetic Mirror**





$$B_x = -\frac{2Cxz^3}{a^4}$$

$$B_y = -\frac{2Cyz^3}{a^4}$$

$$Bz = C\left(1 - \frac{z^4}{a^4}\right)$$

Gyrokinetic + Phase and Time Parallel method are the most efficient

Simple blending could be more accurate as the problem approaches stead state

<sup>\*</sup>Submitted to Journal



## **Summary**



#### Collisional Radiative

- Boltzmann grouping improves group representation over applied uniform distribution
- Minimization of error, which is expected to accumulate quickly during highly-transient durations, and acceleration of method makes Boltzmann reduction a strong case for future coupled simulations
- Adaptive integration allowed for faster simulations at larger electron temperatures
- Sensitive to selected groups
- Grouping reduces stiffness
- Robust solver capable of handling, Te from 1 1,000 eV and  $10^{16}-10^{24}$  particles
- Initial work on non-Maxwellian CR has begun and early numerical issues have been addressed

### Hybridization

 Shows that hybridization technique can be more computational efficient than the schemes that comprise it



## **Future Work**



#### Collisional Radiative Simulations

- Further comparisons between reduced mechanisms with QSS
- Line identification and width assignments in conjunction with experimental spectra
- Apply Boltzmann grouping to new CR Argon rates and test with 1D MHD Argon shock by Kapper, et.al. Future plans to extend to 2 and 3D MF
- Rerun previous LPI test case with level grouping; laser source term, heavy-electron elastic collisions, and multi-electron energy correction and heavy energy equation (Te, Th)
- Group selection through numerical analysis, optimization
- Hybridize Maxwellian CR with QSS
- Non-Maxwellian finite rate EEDF implementation





# Questions?